

Higgs Lineshape

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SNU

Higgs Mini-Workshop: Precision, Cosmology and BSM

with Jeonghyeon Song, Yeo Woong Yoon

1505.00291, 1510.03450, 1601.00006

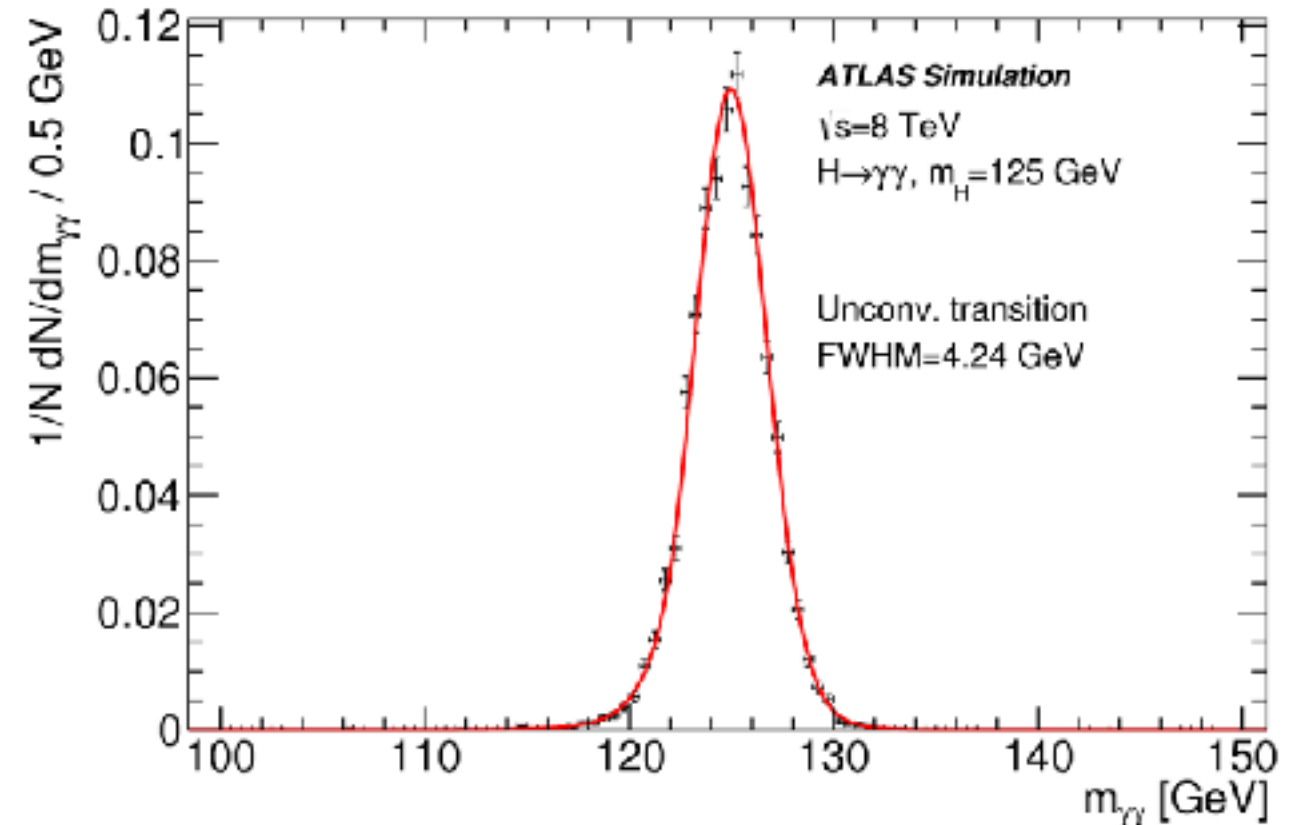
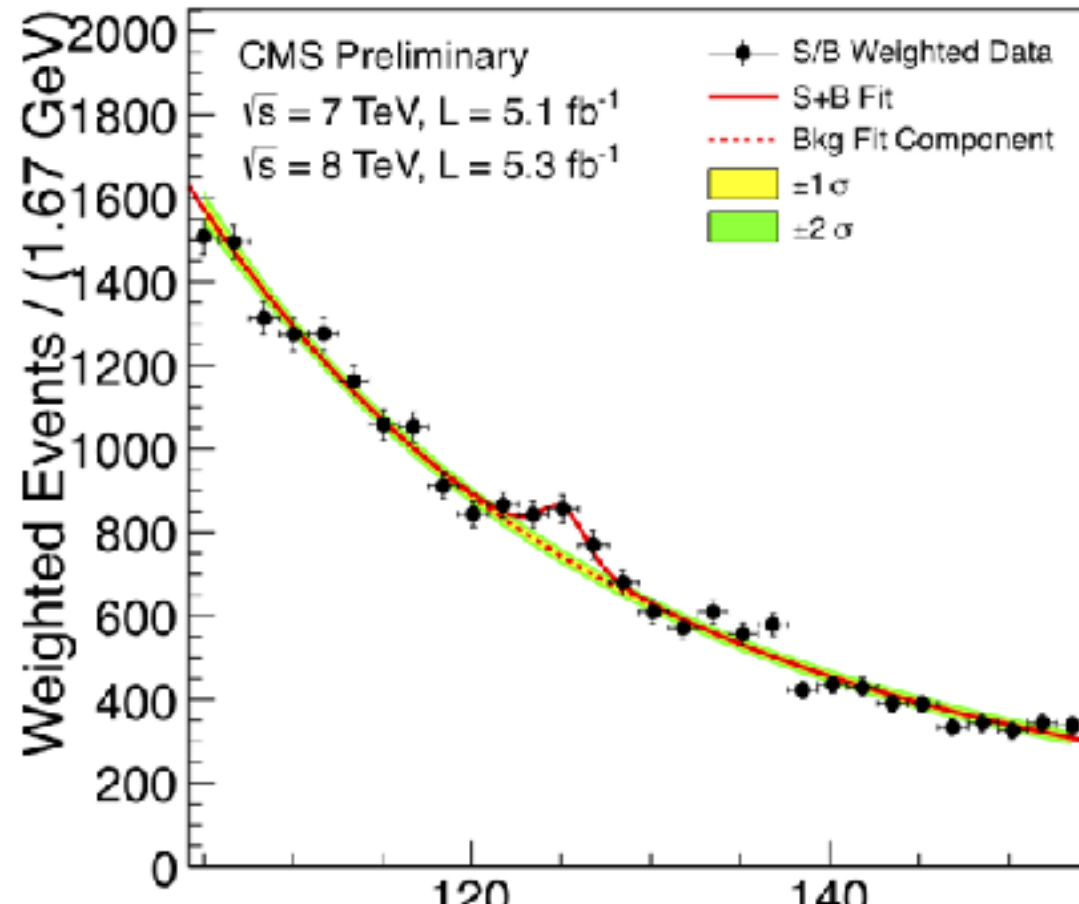
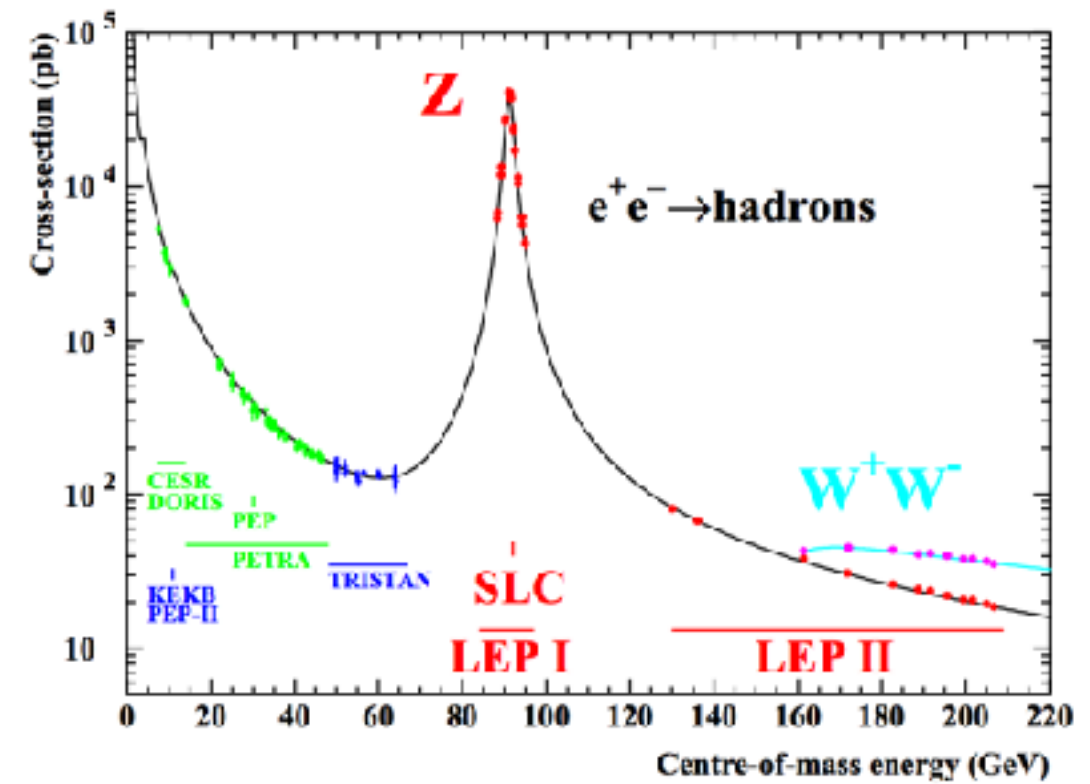
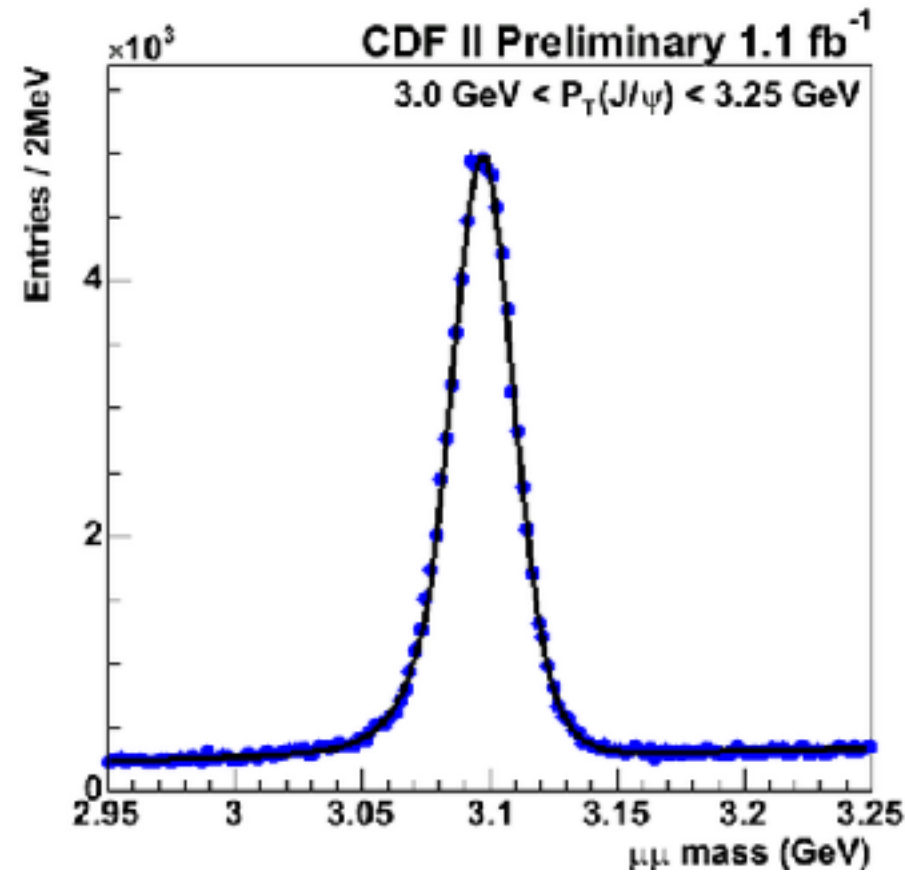
Overlooked Higgs Physics: Resonance Dips

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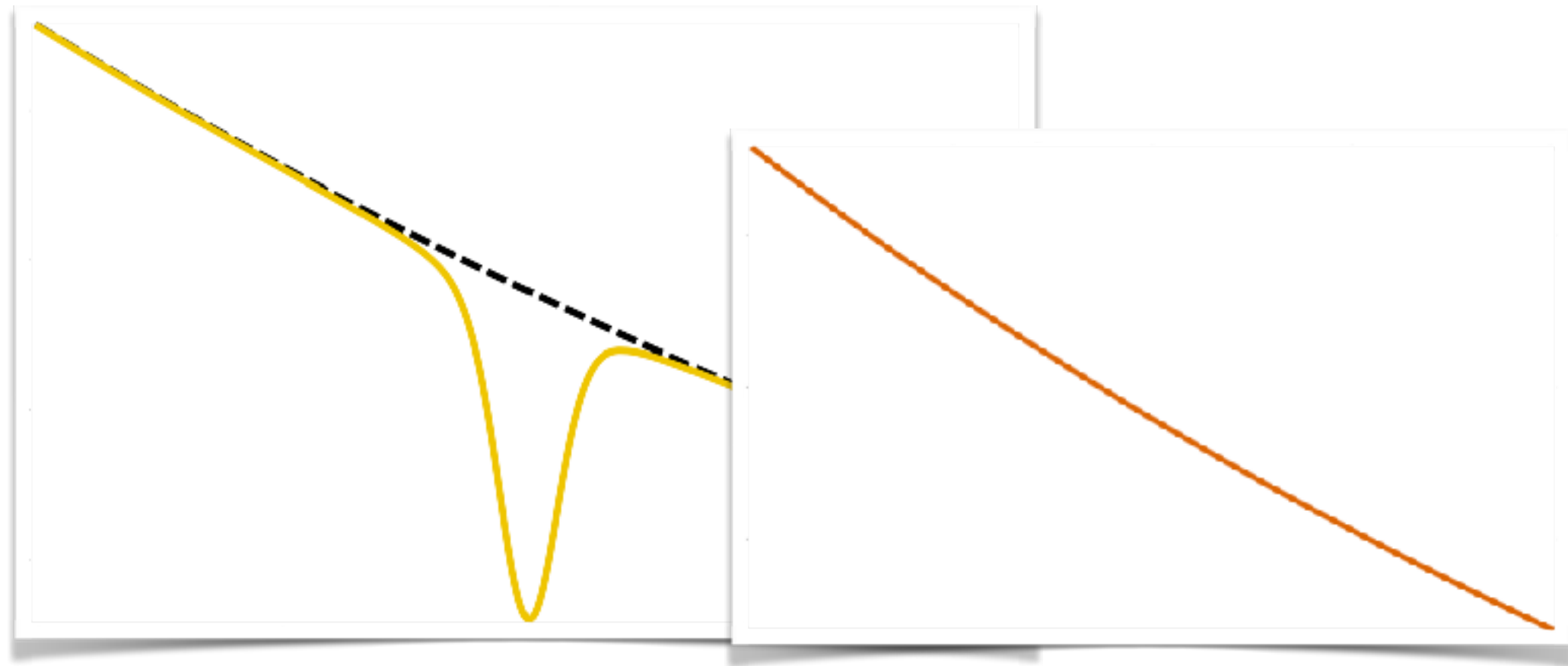
Higgs Mini-Workshop: Precision, Cosmology and BSM

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Resonances in history



A Resonance (pure) Dip



- Good reasons why not seen in elementary particle physics. Also why will be more generic phenomena.
- Not just a mere mathematical possibility.

What my talk is about

1. The resonance dip is a qualitatively new interference phenomena.
 - what is it? and why now?
2. a generic phenomena in 2HDM and MSSM.
 - which channels? what parameters?
 - general rule of thumb
3. inducing new LHC challenges, but also exciting possibilities

1. The resonance dip is a qualitatively new interference phenomena.

- what is it? and why now?

Resonance–continuum interference

- Breit-Wigner resonance peak can be modified from the interference with continuum.

$$\frac{d\hat{\sigma}}{dz} = \frac{1}{32\pi\hat{s}} \sum \left| \mathcal{M}_{\text{bg}} + \frac{M^2}{\hat{s} - M^2 + iM\Gamma} \cdot \mathcal{M}_{\text{res}} \right|^2$$

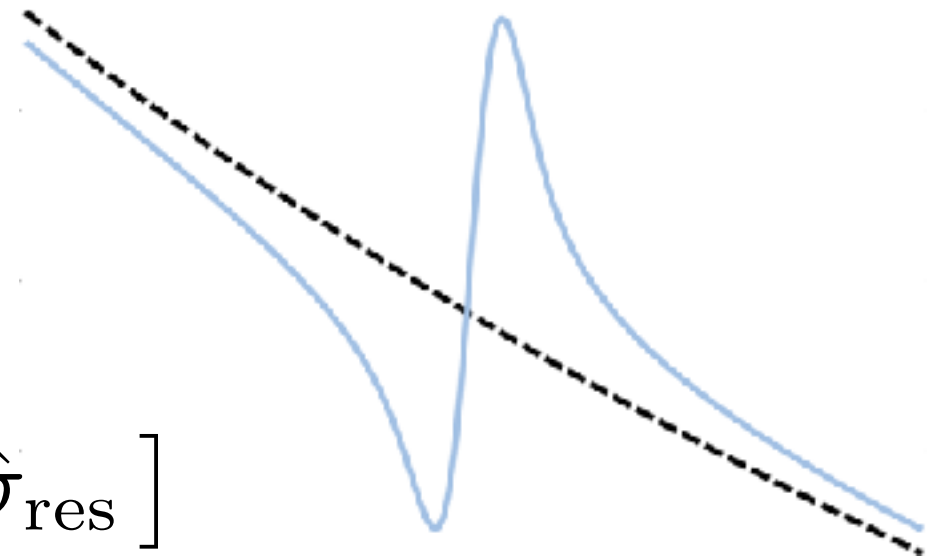
Purely real interference

- Breit-Wigner resonance peak can be modified from the interference with continuum.

$$\frac{d\hat{\sigma}}{dz} = \frac{1}{32\pi\hat{s}} \sum \left| \mathcal{M}_{\text{bg}} + \frac{M^2}{\hat{s} - M^2 + iM\Gamma} \cdot \mathcal{M}_{\text{res}} \right|^2$$

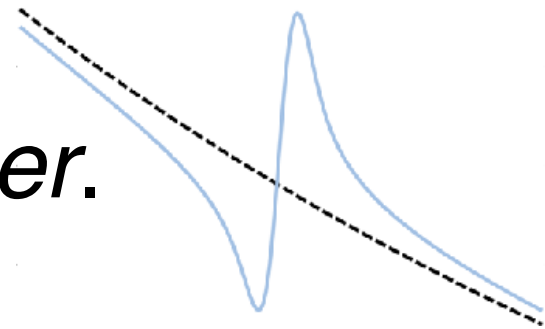
- A typical example: **purely real intf**

$$\hat{\sigma} = \hat{\sigma}_{\text{bg}} + \frac{1}{(\hat{s} - M^2) + M^4 w^2} \times \left[M^2(\hat{s} - M^2)\hat{\sigma}_{\text{int}} + M^4\hat{\sigma}_{\text{res}} \right]$$

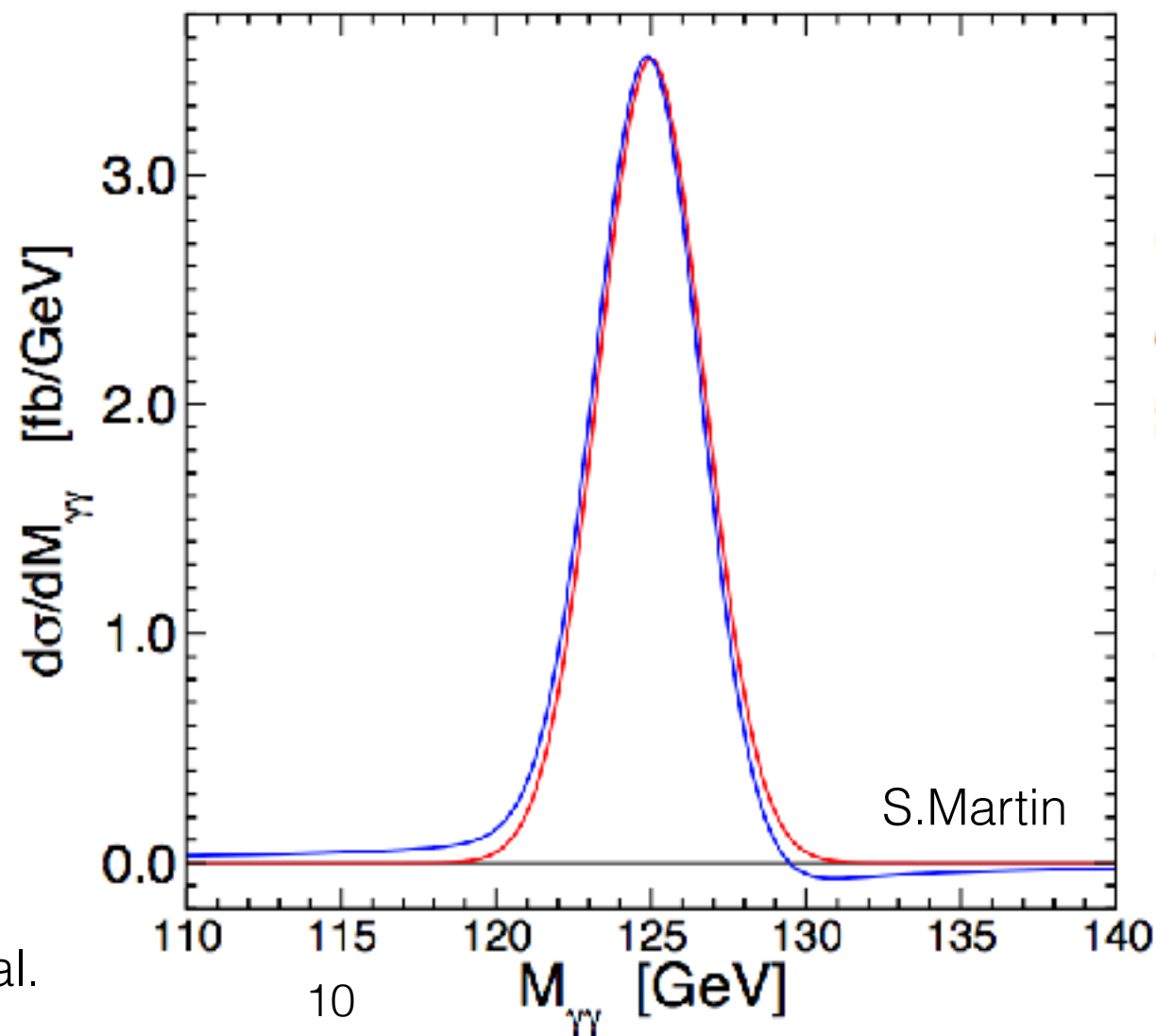


Purely real interference

- The **real-part** induces a peak and dip *together*.
- Shifting the location of a peak (affecting the pole mass measurement).



~70 MeV shift to the
125 GeV Higgs mass

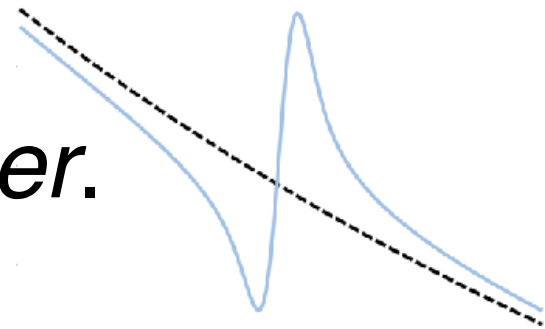


L.Dixon et al.

Sunghoon Jung

Purely real interference

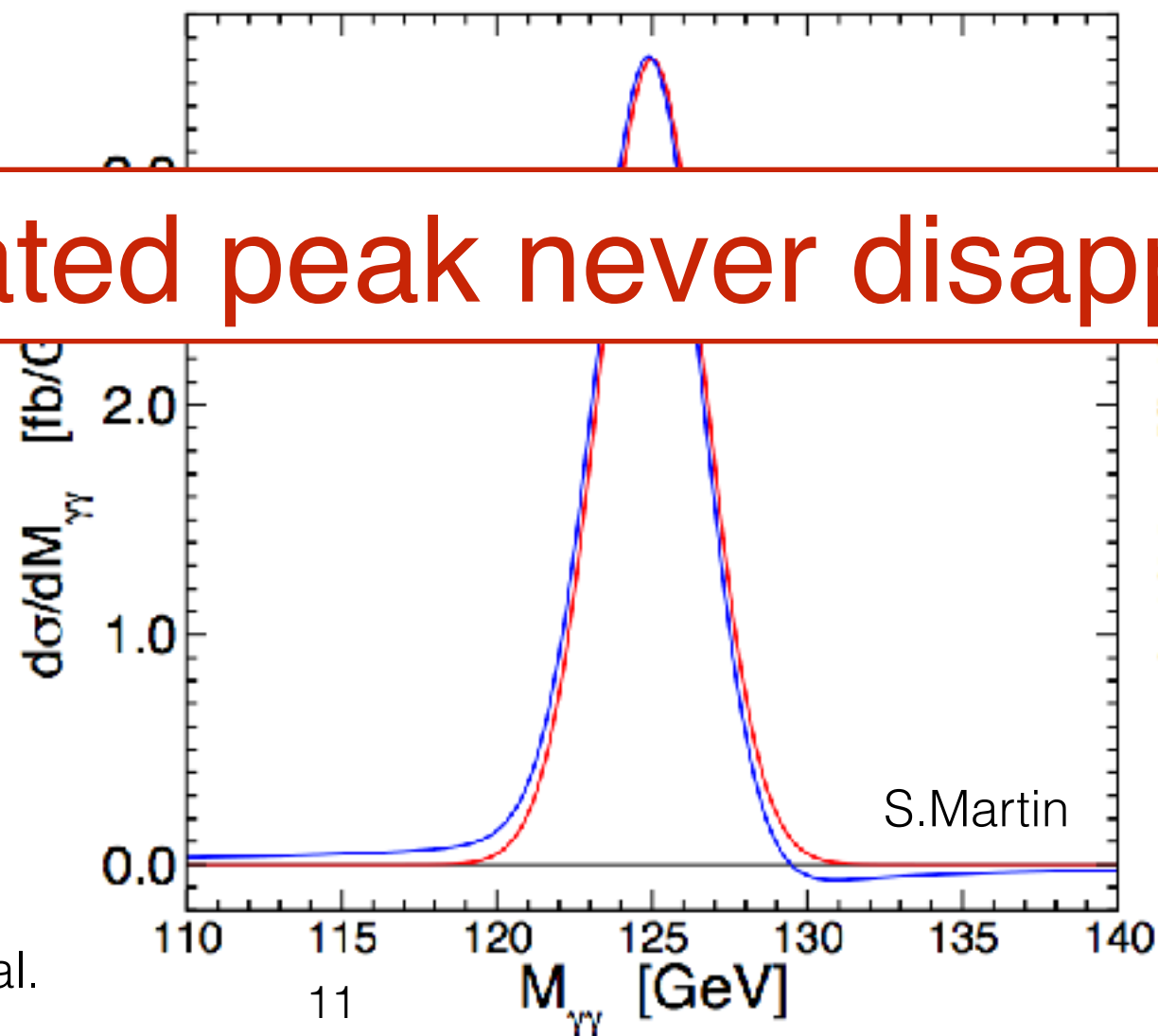
- The **real-part** induces a peak and dip *together*.
- Shifting the location of a peak (affecting the pole mass measurement).



But the associated peak never disappears!

~70 MeV shift to the
125 GeV Higgs mass

L.Dixon et al.



Now with a relative phase

SJ, Song, Yoon

- The resonance-continuum interference **with a relative phase.**

$$\frac{d\hat{\sigma}}{dz} = \frac{1}{32\pi\hat{s}} \sum \left| A_{\text{bg}} e^{i\phi_{\text{bg}}} + \frac{M^2}{\hat{s} - M^2 + iM\Gamma} \cdot A_{\text{res}} e^{i\phi_{\text{res}}} \right|^2$$

Now with a relative phase

SJ, Song, Yoon

- The resonance-continuum interference **with a relative phase.**

$$\frac{d\hat{\sigma}}{dz} = \frac{1}{32\pi\hat{s}} \sum \left| A_{\text{bg}} e^{i\phi_{\text{bg}}} + \frac{M^2}{\hat{s} - M^2 + iM\Gamma} \cdot A_{\text{res}} e^{i\phi_{\text{res}}} \right|^2$$

$$\hat{\sigma} = \hat{\sigma}_{\text{bg}} + \frac{1}{(\hat{s} - M^2) + M^4 w^2}$$

$$\times \left[M^2(\hat{s} - M^2) \hat{\sigma}_{\text{int}} c_\phi + M^4 \hat{\sigma}_{\text{res}} \left(1 + \frac{2w}{R} s_\phi \right) \right]$$

$$w = \Gamma/M$$

$$R = \frac{\hat{\sigma}_{\text{res}}}{\hat{\sigma}_{\text{int}}} \simeq \frac{A_{\text{res}}}{A_{\text{bg}}}$$

$$\phi \simeq \phi_{\text{res}} - \phi_{\text{bg}}$$

typical real-part
interference

resonance
square

imaginary
interference

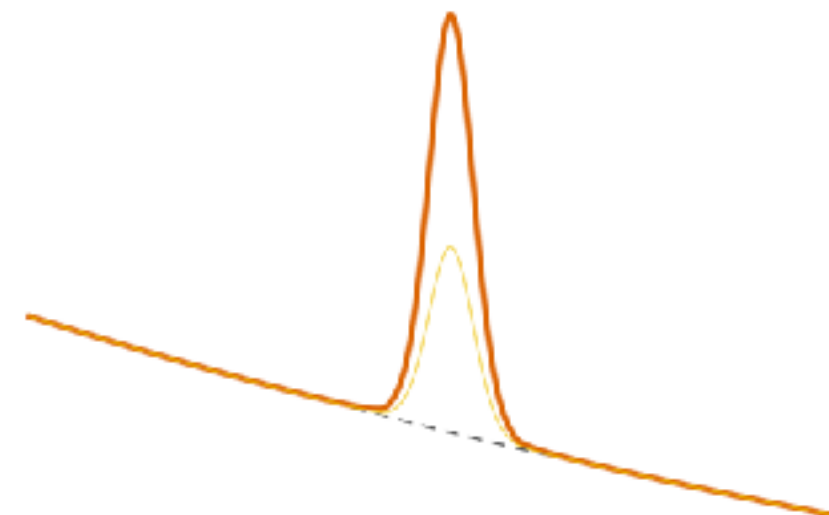
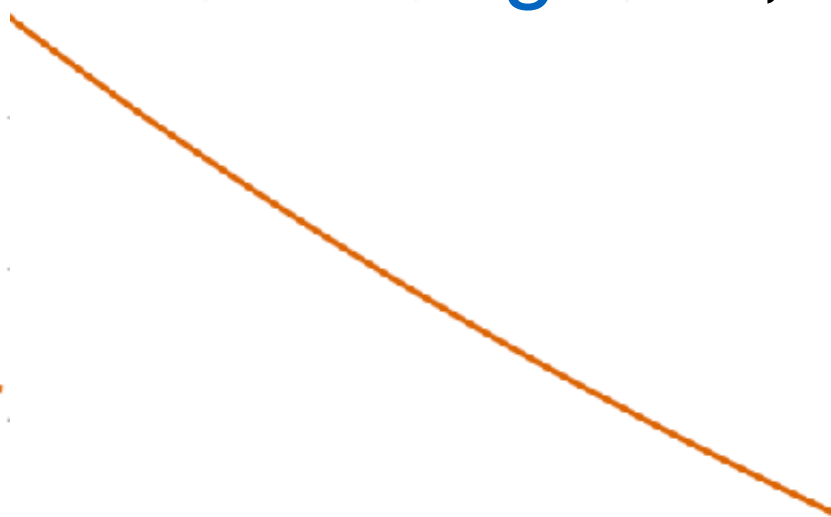
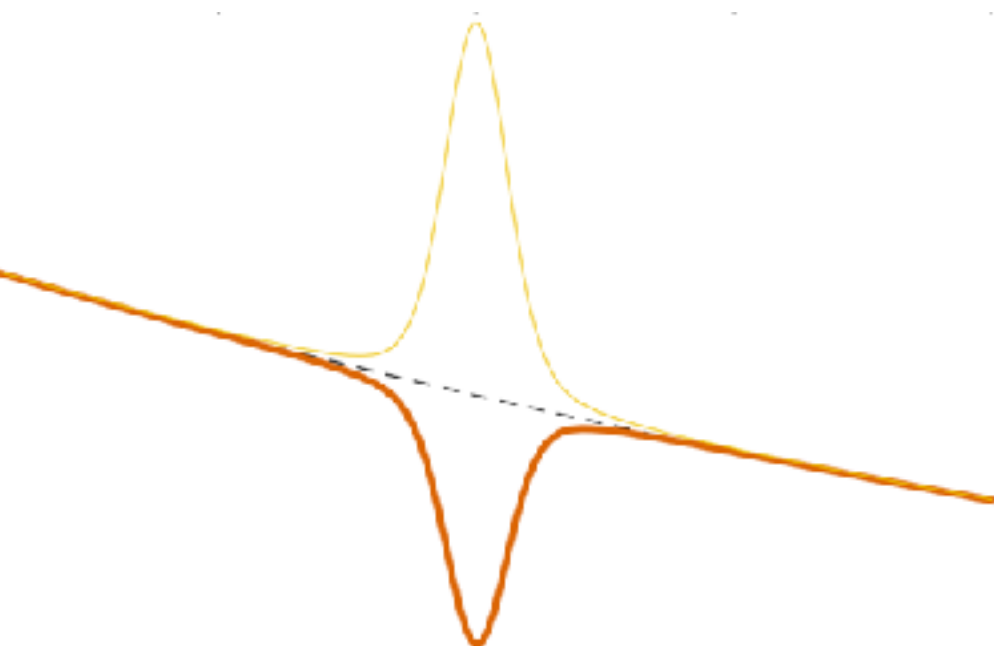
Purely imaginary intf

SJ, Song, Yoon

- New lineshapes:
pure dip, nothingness, and rescaled pure peak.

$$\hat{\sigma} = \hat{\sigma}_{\text{bg}} + \frac{1}{(\hat{s} - M^2) + M^4 w^2} \left[0 + M^4 \hat{\sigma}_{\text{res}} \left(1 \pm \frac{2w}{R} \right) \right]$$

negative, vanishing, positive

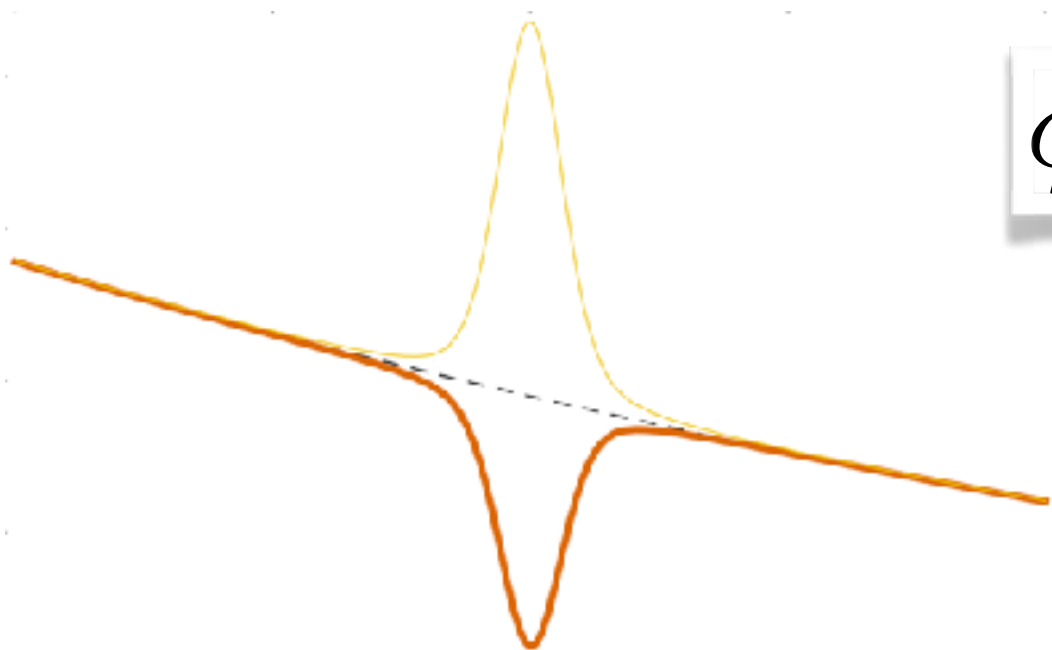


Pure resonance dip

SJ, Song, Yoon

- I'll focus on the **pure dip** in this talk—not only most extreme case, but also turns out to be useful.

$$\hat{\sigma} = \hat{\sigma}_{\text{bg}} + \frac{1}{(\hat{s} - M^2) + M^4 w^2} \left[0 + M^4 \hat{\sigma}_{\text{res}} \left(1 \pm \frac{2w}{R} \right) \right]$$



$$\phi = -\pi/2, \quad w/R > 0.5$$

$$R = \frac{\hat{\sigma}_{\text{res}}}{\hat{\sigma}_{\text{int}}} \simeq \frac{A_{\text{res}}}{A_{\text{bg}}}$$

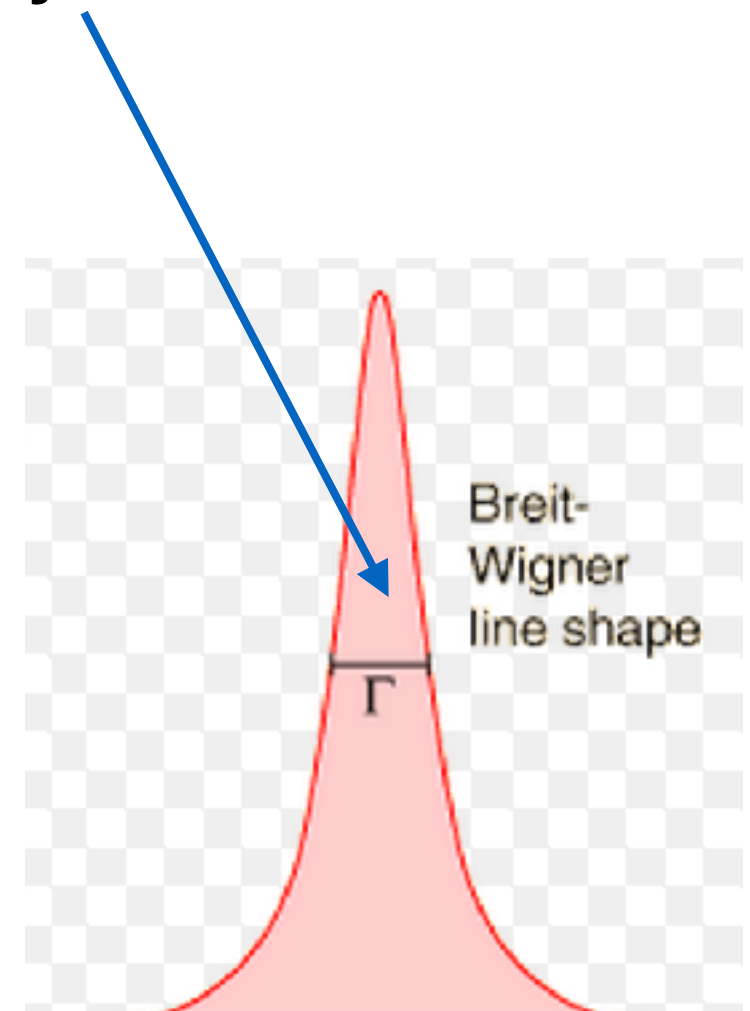
$$w = \Gamma/M$$

Narrow width approximation (NWA)

- The **total resonance rate** is the area below the BW peak, and is estimated approximately by NWA.

$$\frac{1}{(\hat{s} - M^2)^2 + M^2\Gamma^2} \rightarrow \frac{\pi}{M\Gamma} \delta(\hat{s} - M^2)$$

$$\sigma(H/A) \cdot BR(H/A \rightarrow t\bar{t})$$

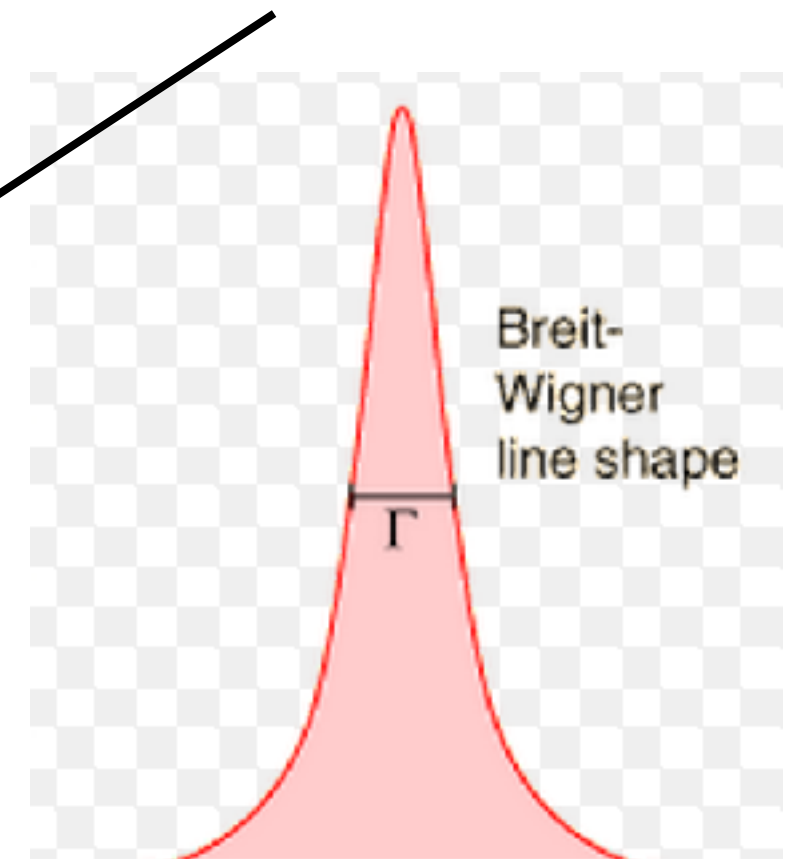


Modified NWA

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- The imaginary-part now does contribute to the total rate; usual NWA is not valid with the imaginary-part.
- We can modify the NWA by the multiplicative C-factor.

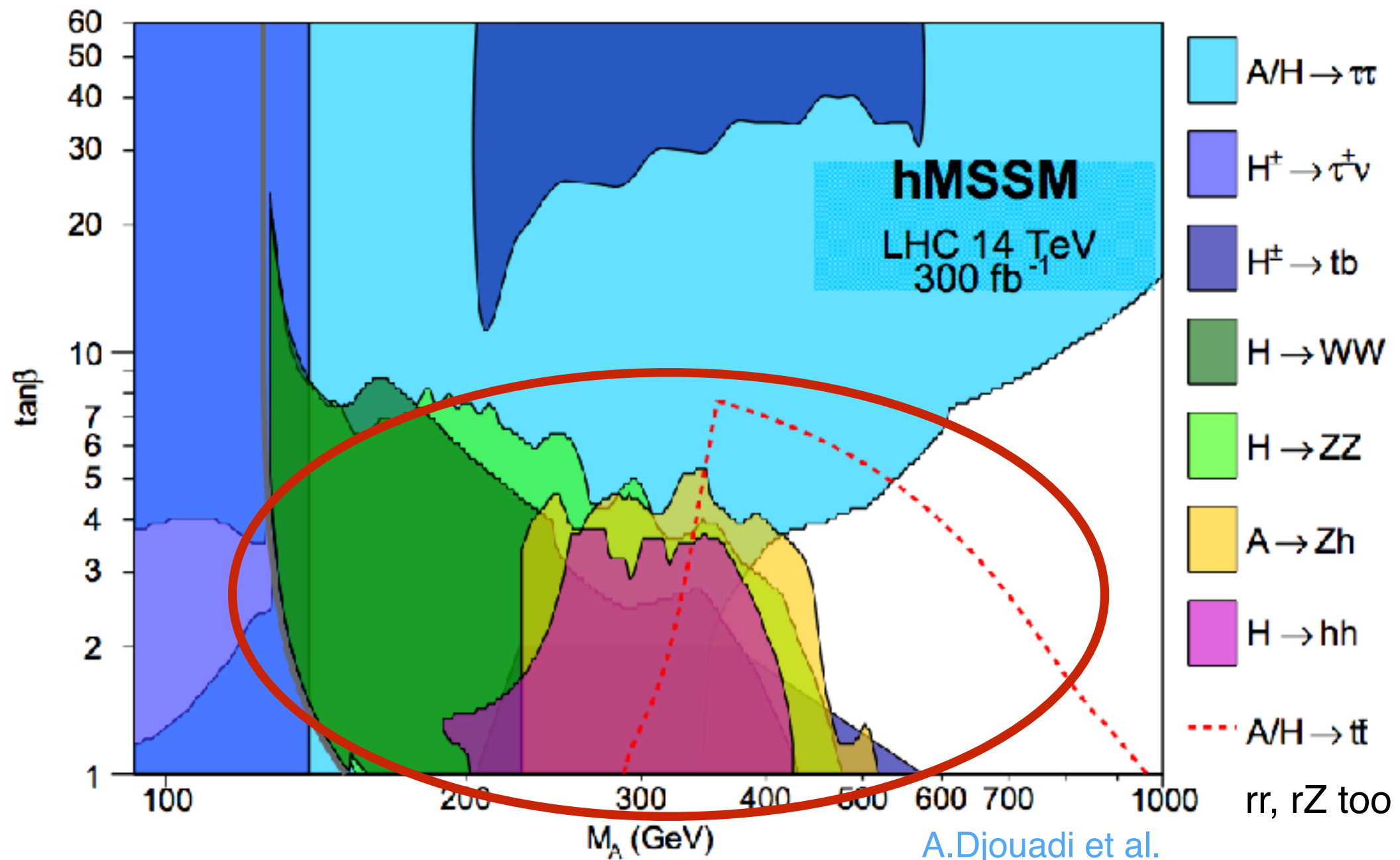
$$\sigma(H/A) \cdot BR(H/A \rightarrow t\bar{t}) \cdot \left(1 + 2\frac{w}{R}s_\phi\right)$$



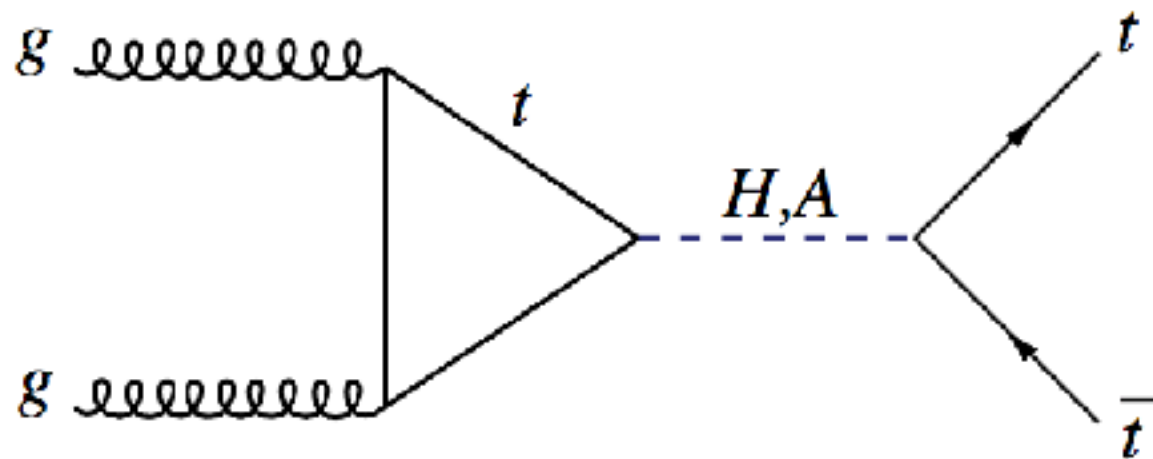
2. It is a generic phenomena in 2HDM and MSSM.

- which channels? what parameters?
- general rule of thumb

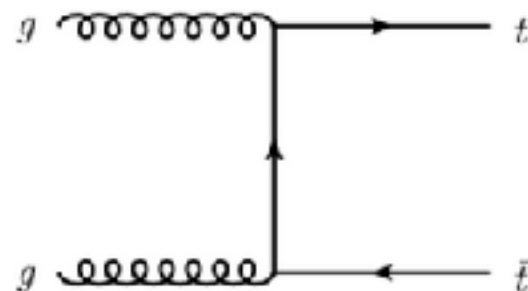
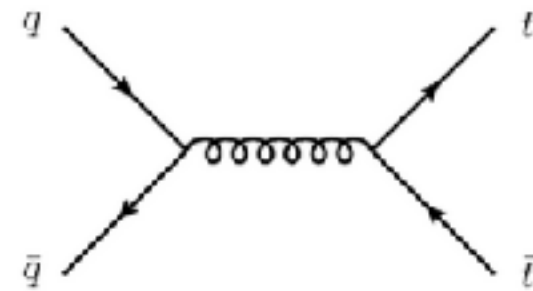
MSSM Heavy Higgs in most channels can be dips!



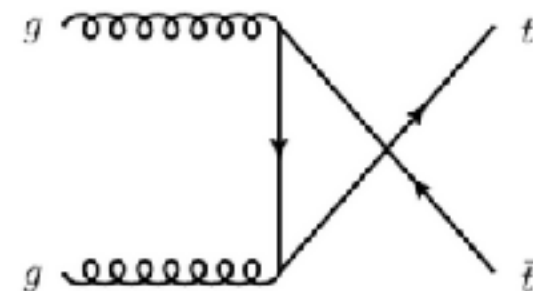
Heavy Higgs in $t\bar{t}$



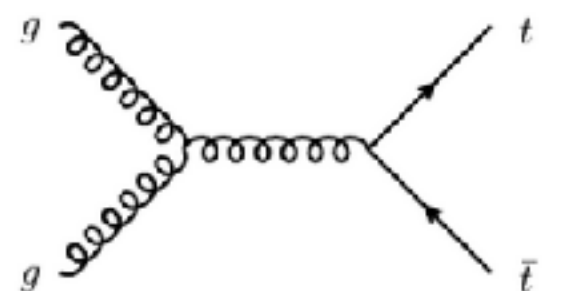
alignment scenario,
but similar in general



+

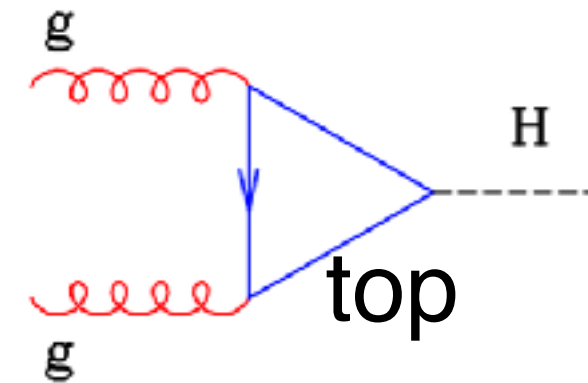


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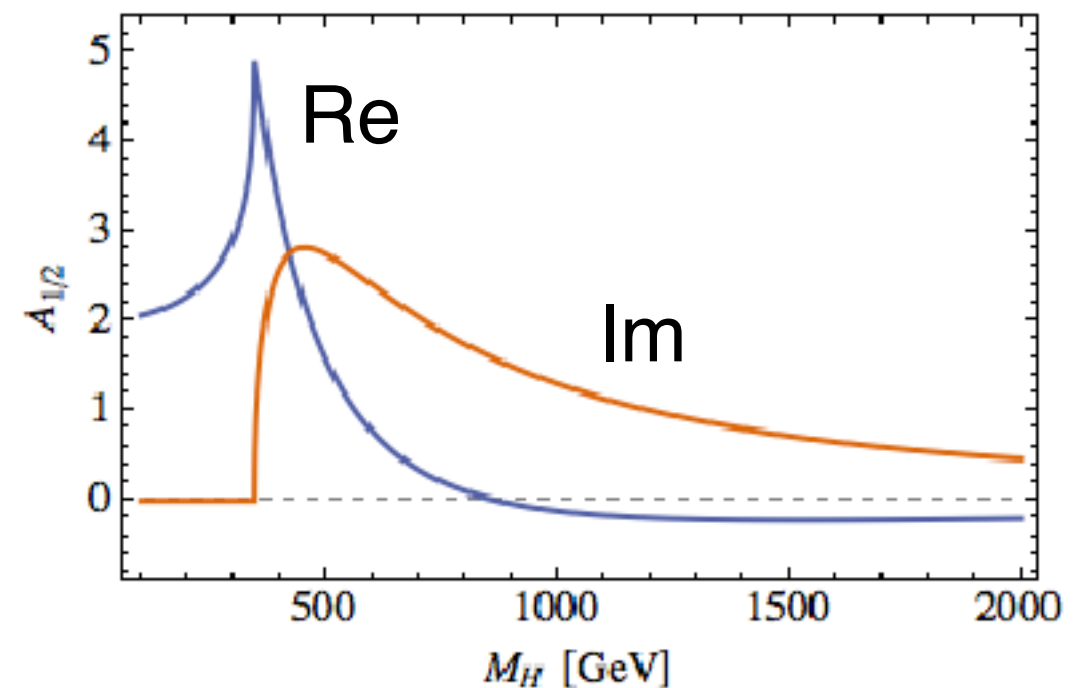


Mechanism 1: vanishing Re.

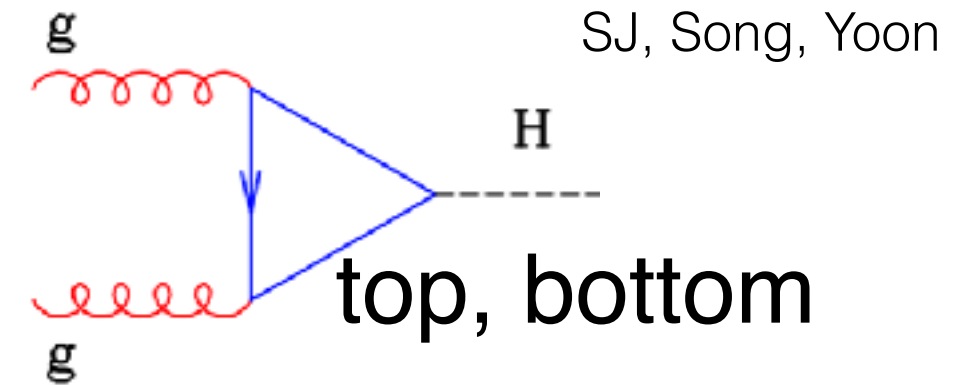
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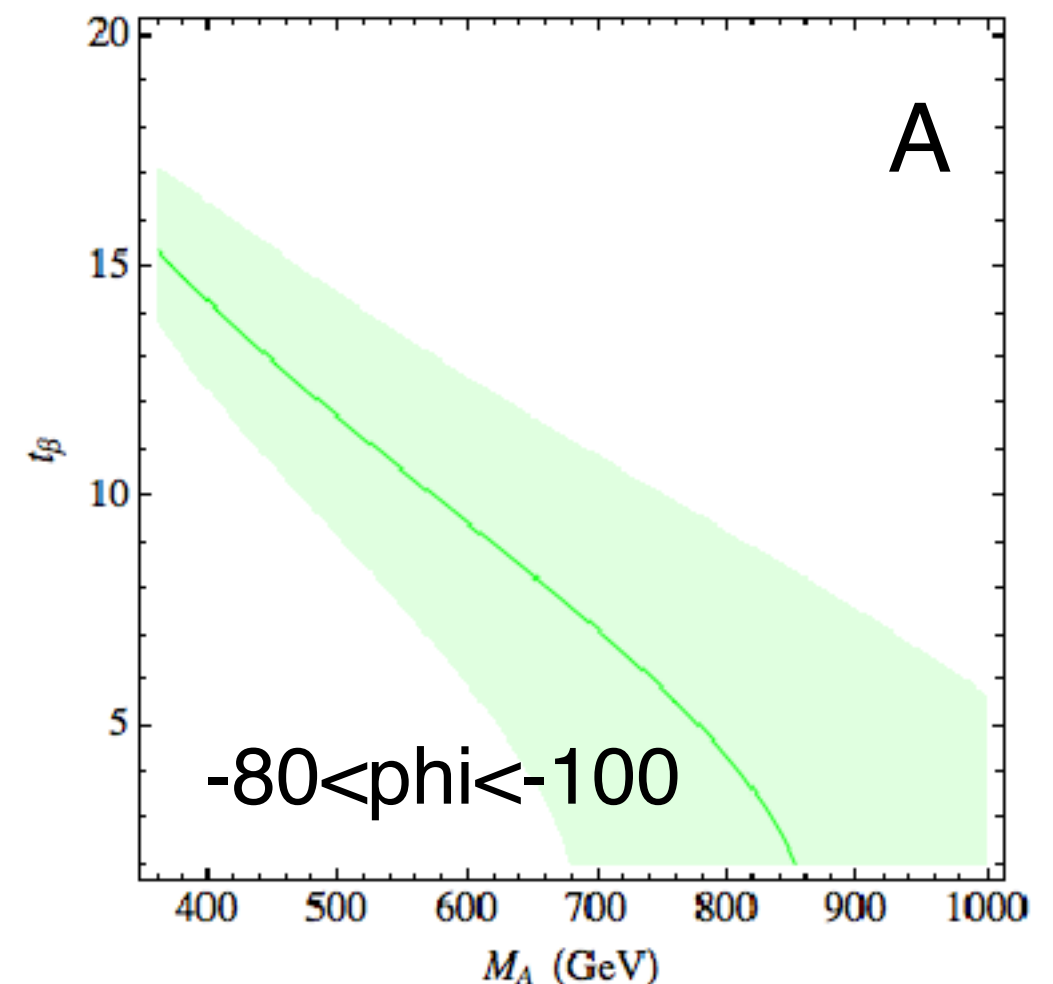
- The triangle loop function can accidentally have a *vanishing* Real part.
- For $t_b=2$, $MA=850$, $MH=1150$ are pure dips.



Mechanism 2: cancelling Re.

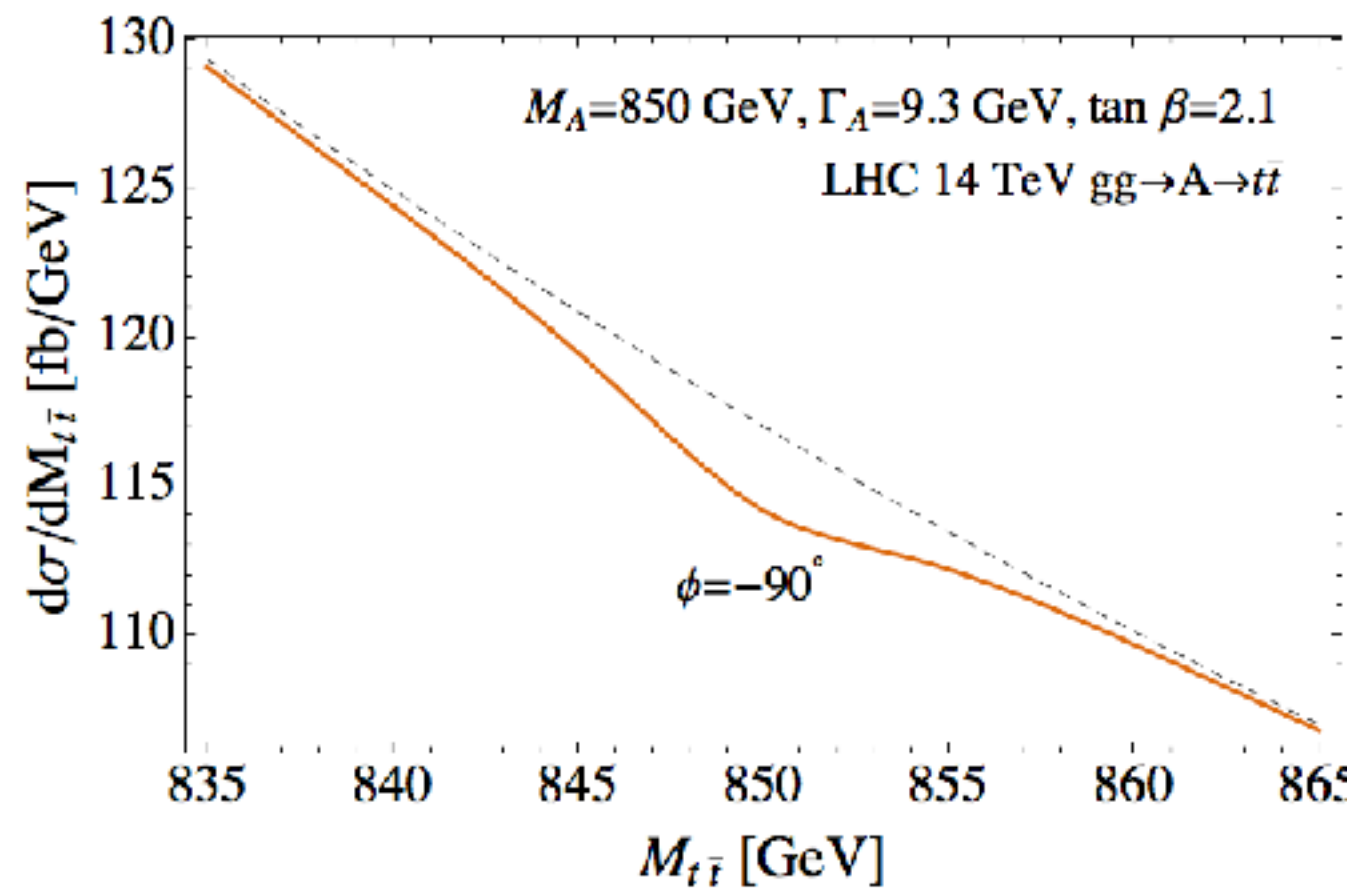
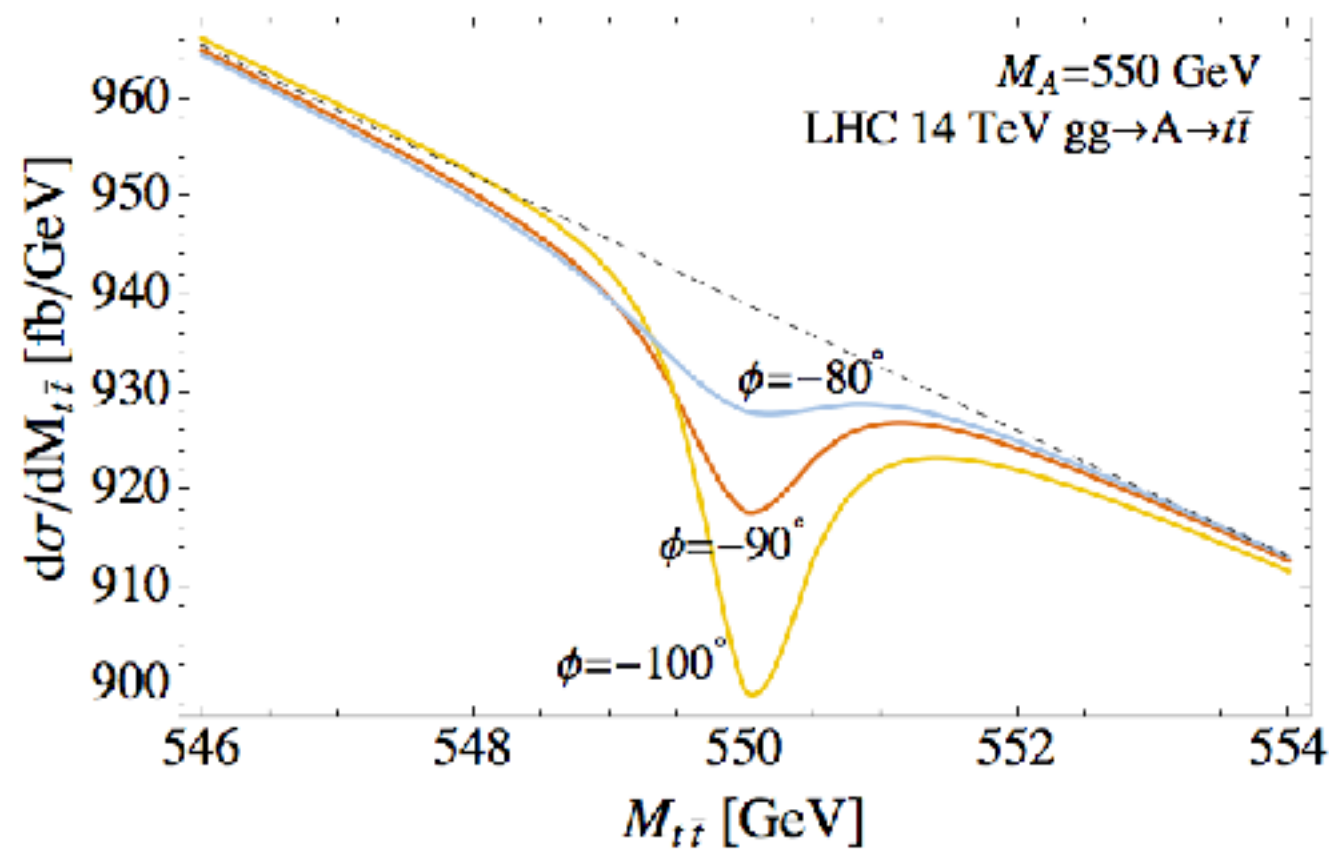


- For larger $\tan \beta$, the t-loop and b-loop can *cancel* their Real parts.
- Any $2m_t < M_A < 850$ can be a pure $t\bar{t}$ dip.



Demonstration

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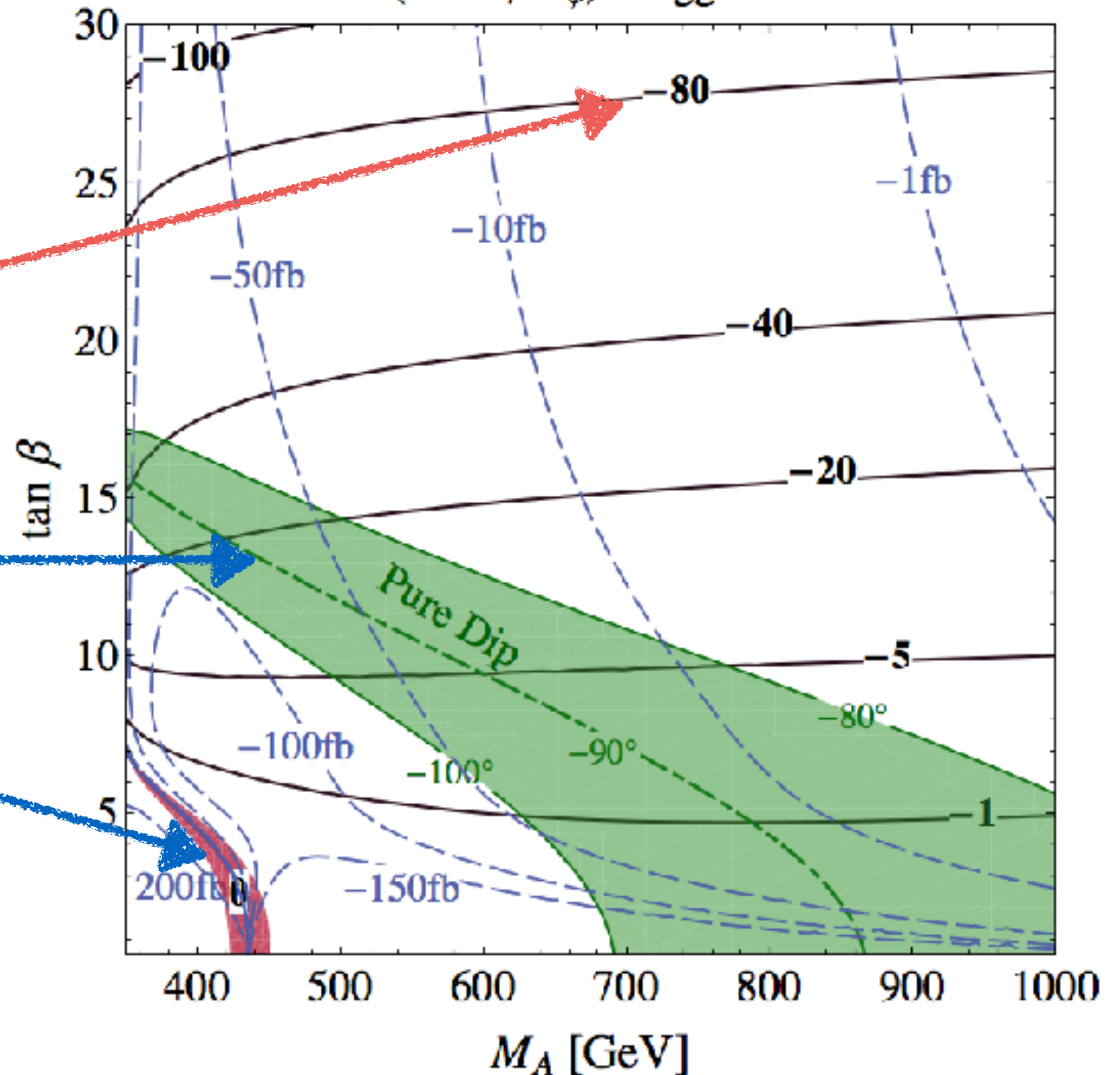


ttbar resonance shapes

$C = (1 + 2w/Rs_\phi)$ for $gg \rightarrow A \rightarrow t\bar{t}$ SJ, Song, Yoon

$$C = \left(1 + 2\frac{w}{R}s_\phi\right)$$

- Pure dip and nothingness regions.

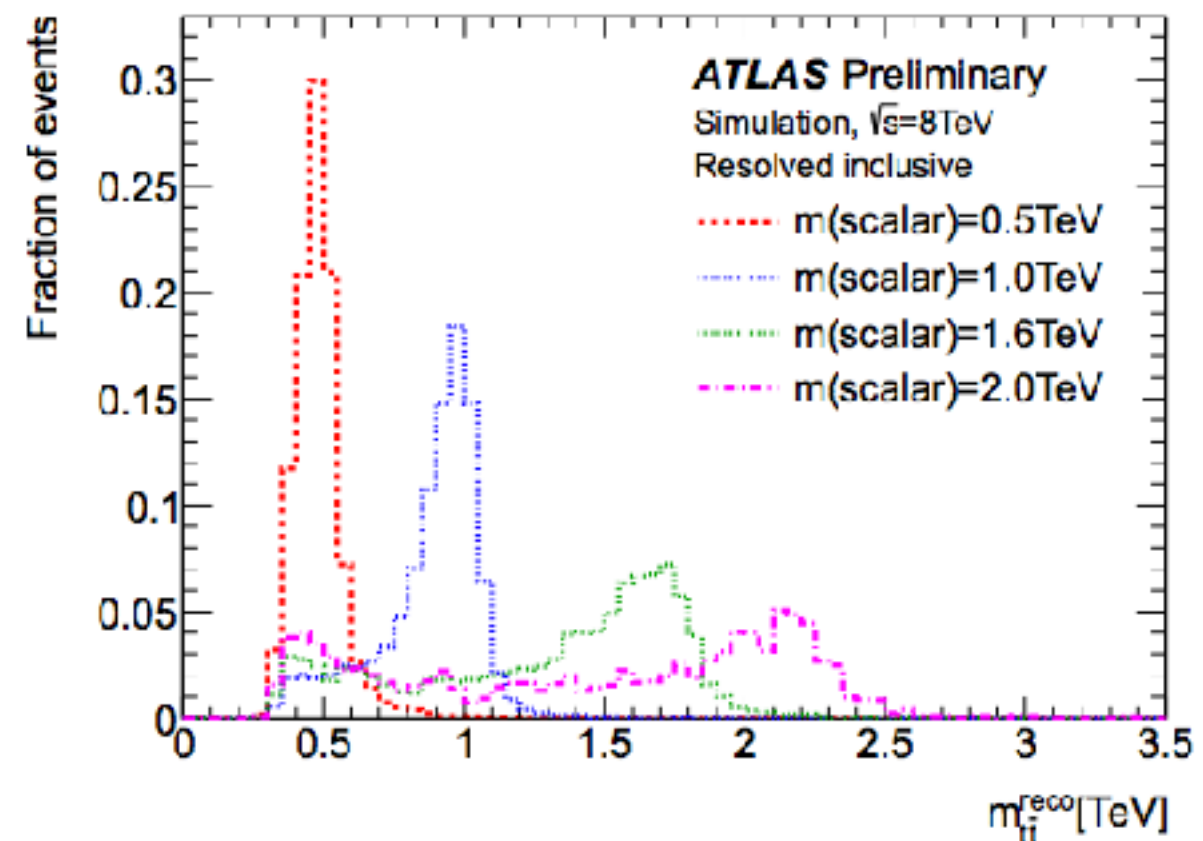


3. New LHC challenges and exciting possibilities

LHC Challenges

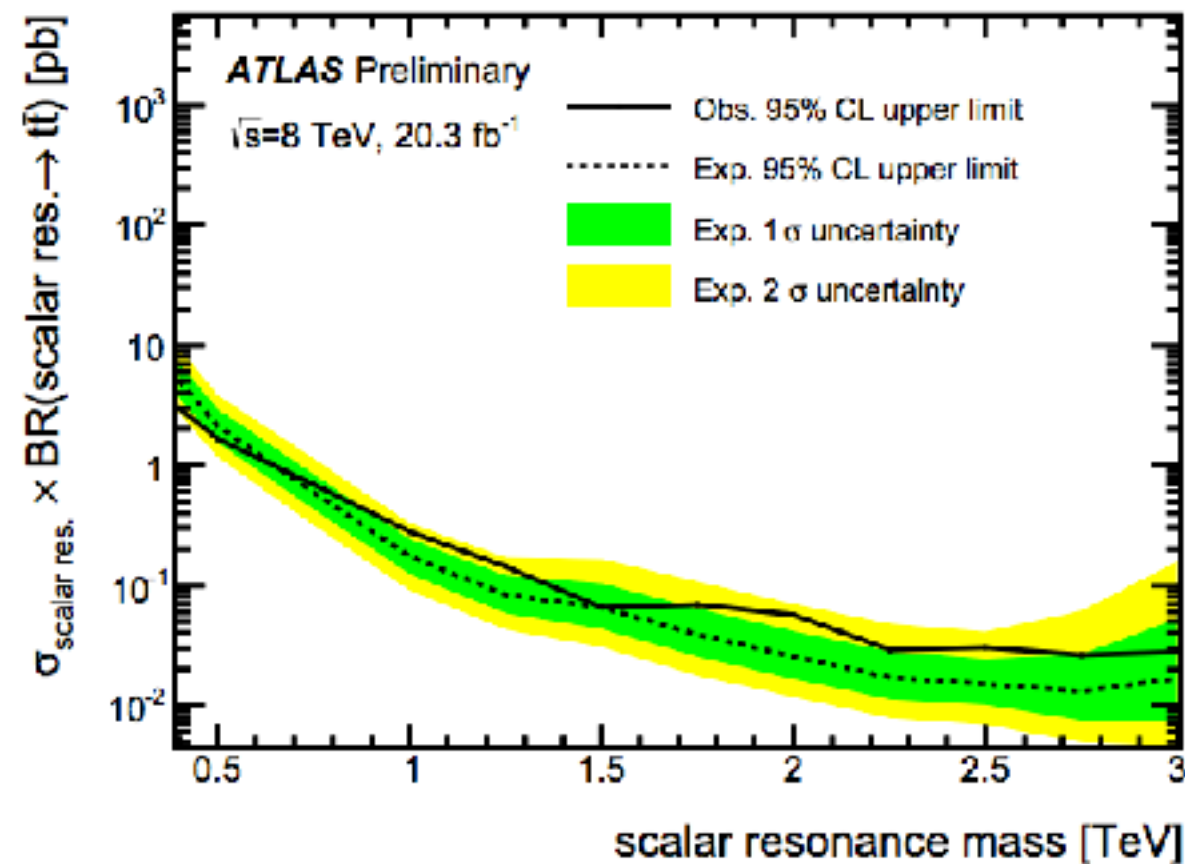
- **M(ttbar) resolution** is $O(10)$ GeV at best, smearing/ canceling out any delicate shapes.
- **Huge simulation time** (varieties of line shapes, loop processes).

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(CMS: 1309.2030)



NB: one can look for other channels: ttH, bbH, etc.
(J.Hajer et al, N.Craig et al.)

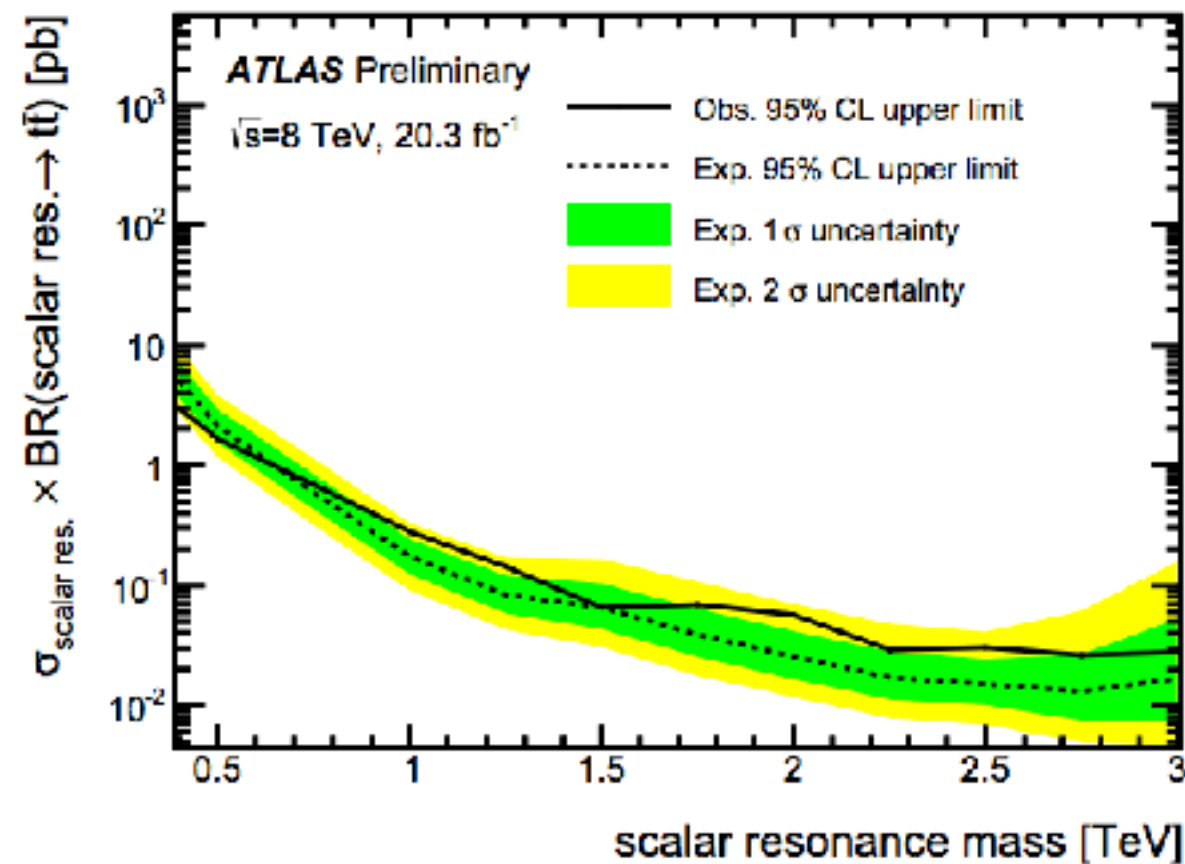
Current H/A- $\rightarrow$ ttbar search



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(CMS: 1309.2030)

- Currently, no interferences are accounted for. (but all realistic effects are accounted)
- No predictions overlaid b/c no such scalar model.

Current H/A- $\rightarrow$ ttbar search



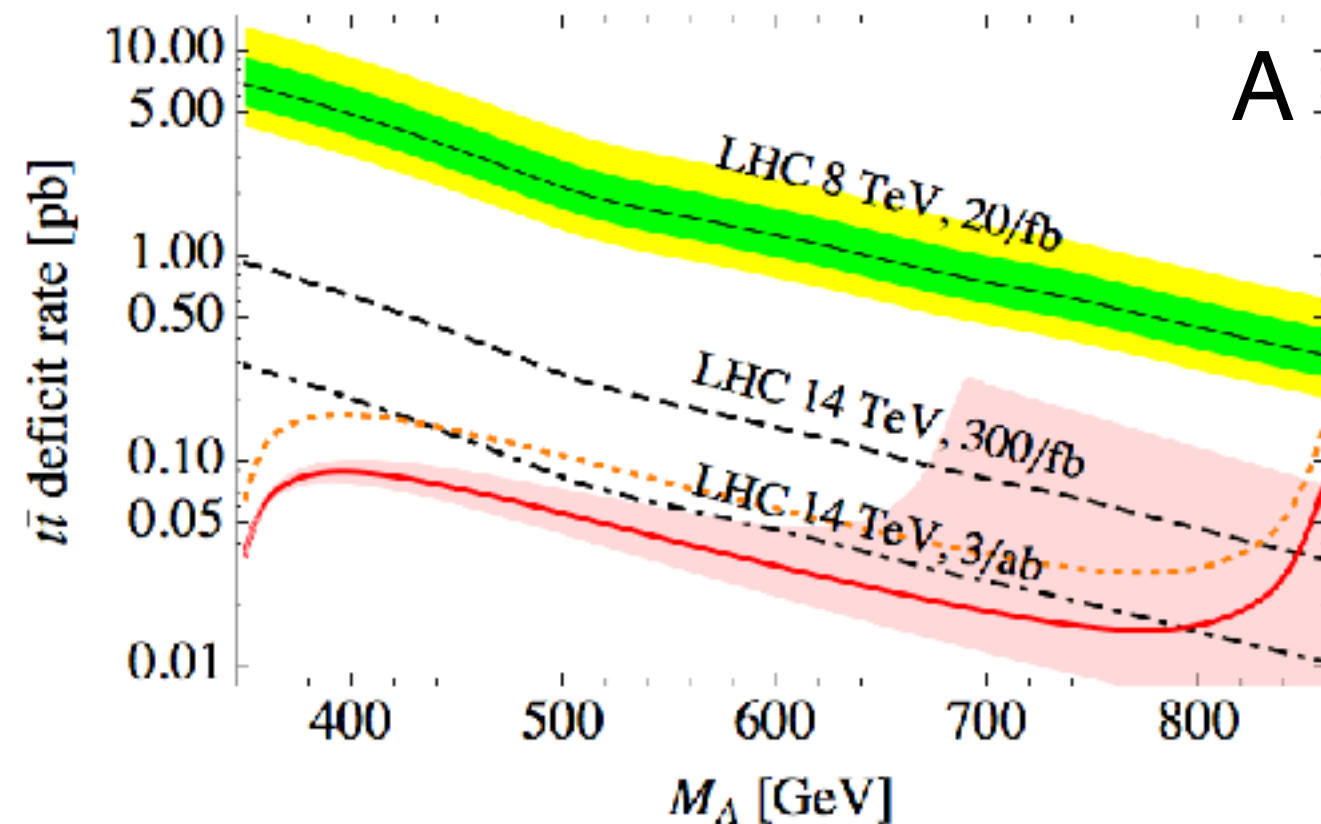
ATLAS-CONF-2015-009
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BUT, this exp result is applicable to pure dips
(even w/o interference accounted for)!

Sensitivities to pure dip

- Same shapes for BW peaks and pure dips **as if no interferences.**
- Searches inherently sensitive to both **excess(peak)** and **deficit(dip)**—just deviation.
- Thus, current results apply to pure dips **as is now.**

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$$\sigma(H/A) \cdot BR(H/A \rightarrow t\bar{t}) \cdot \left(1 + 2\frac{w}{R}s_\phi\right)$$

How Resonance-Continuum Interference Changes 750 GeV Diphoton Excess: Signal Enhancement and Peak Shift

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Summary

- The qualitatively new interference effect generally arise in **MSSM/2HDM** in heavy Higgs Golden channels.
- **A few ideas to circumvent the LHC challenges** were proposed, but many more are plausible/needed.
- Can arise more generically in **various new models and channels**.

Are we missing any deficits or nothingness?!